

Non-Stationary Processes in Atmospheres of Early-Type Stars: Influence on Forbidden to Intercombination Ratio f/i

V.V. Dushin¹, A.F. Kholtygin¹

E-mail: *v.dushin@spbu.ru*

We report the results of non-stationary level population modeling of highly ionized atoms in the atmospheres of early-type stars. We studied the influence of the fast heating and cooling processes on the ratio of forbidden to intercombination line intensities $R = f/i$ for He-like ions (CV, NVI, OVII, etc.) in X-ray spectra.

It is shown that the instantaneous ratio R_m for the non-stationary plasma varies by up to 4 orders of magnitude on short time scales (milliseconds) in comparison with the value for the stationary plasma. In the same time the value of R_a averaged on long time scales (hours and minutes) varies by up to 20%. Using the ratio R calculated in the case the stationary plasma for the non-stationary plasma can lead to an overestimation of the plasma electron density by up to 1–2 orders of magnitude.

1 Introduction

The density diagnostics of the X-ray emitted plasma of the early-type stars based on the forbidden-to-intercombination line ratio R showed that this ratio is much lower than that predicted for the homogeneous non-stationary plasma (e.g., [1]). This could be explained as follows:

- The stellar UV-radiation excites electrons in the upper level $1s2s\ ^3S_1$ of the forbidden line f and populates the upper level $1s2p\ ^3P_{1,2}$ of the intercombination line i , which weakens the line f and strengthens the line i .
- Bound electrons are excited from the upper level of the line f by free electrons, which decreases the line f . This happens if the X-ray radiation originates from the dense clouds in stellar atmospheres.

In this paper we outline an alternative hypothesis.

2 Processes in non-stationary plasma

We suppose that the level population in the stellar atmospheres can become non-stationary.

¹ St. Petersburg State University, Russia

Non-stationarity could be caused by collisions of the plasma flows in the region on the stellar magnetic equator or by nano-flares in the stellar atmosphere similarly to the solar ones.

The non-stationary level population can be described by the following equation:

$$\frac{dx_i}{dt} = n_e \sum_{j \neq i}^N x_j q_{ji} + \sum_{j=i+1}^N x_j A_{ji} - x_i \left(\sum_{j=1}^{i-1} A_{ij} + n_e \sum_{j \neq i}^N q_{ij} \right). \quad (1)$$

Here x_i is the relative population of the i -th level, N is the number of levels considered, n_e is the electron density, q_{ji} is the excitation/deexcitation rate from the level j to the level i , A_{ij} is the corresponding Einstein A -value. In our models we used $N = 50$ levels, which was enough for a precise modeling.

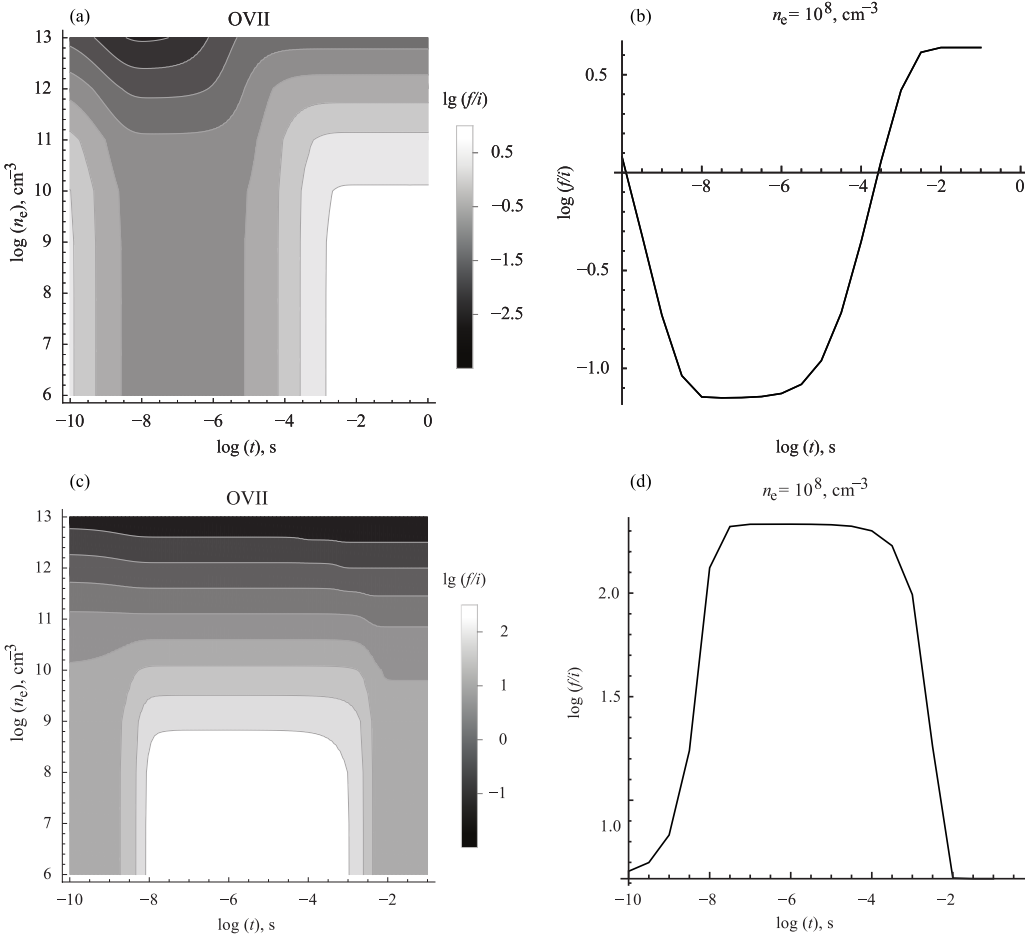


Figure 1: *Panel a*: Dependence of the ratio R_m on time and electron number density for OVII in the model for rapidly heated plasma. At $t = 0$, the plasma instantaneously heats from $T_e = 10^6$ K to $T_e = 10^7$ K, after the heating the temperature remains constant. *Panel b*: the same but for $n_e = 10^8 \text{ cm}^{-3}$. *Panels c, d*: the same as the upper row, but for rapid cooling from $T_e = 10^7$ K to $T_e = 10^6$ K.

Table 1: The parameters of the models: T_c is the plasma temperature in the “cool” state, t_c is the time of cooling, R_c is the stationary line ratio for T_c , the similar parameters for plasma heating are indexed with “h”

Model	T_c , K	t_c , s	T_h , K	t_h , s	R_c	R_h
A	5×10^5	3×10^{-2}	10^8	10^{-5}	9×10^{-3}	4×10^{-2}
B	5×10^5	10^{-5}	10^8	10^{-5}	3	2.75
C	5×10^5	3×10^{-2}	10^8	10^{-5}	9×10^{-4}	4×10^{-3}
D	5×10^5	10^{-3}	10^7	3×10^{-3}	7.5×10^{-1}	1.34
E	5×10^5	3×10^{-2}	10^8	10^{-5}	4.44	2.94
F	5×10^5	3×10^{-2}	10^8	10^{-5}	4.26	2.93
G	5×10^5	3×10^{-3}	10^7	2×10^{-3}	1.78	2.94
H	5×10^5	10^1	10^7	10^{-1}	5.61	8.25
I	5×10^5	10^{-3}	10^8	10^{-5}	2.98	2.75

Table 2: The results of the modeling, where n_e is the input model electron density, R_a is the averaged line ratio f/i , \bar{n}_e is the electron density derived from R_a , when supposing the stationarity of the plasma

Model	Ion	n_e , cm^{-3}	R_a	\bar{n}_e , cm^{-3}	$\lg(n_e/\bar{n}_e)$
A	OVII	10^{13}	10^{-2}	8×10^{12}	0.1
B	OVII	10^{10}	–	–	–
C	OVII	10^{14}	2×10^{-3}	5×10^{13}	0.3
D	OVII	10^{11}	1.18	5×10^{10}	0.3
E	OVII	10^8	3.48	6×10^9	–1.8
F	OVII	10^9	3.35	6×10^9	–0.8
G	NVI	10^{10}	2.12	8×10^9	0.1
H	CV	10^9	6.24	8×10^8	0.1
I	OVII	10^{10}	2.93	–	–

3 Modeling and discussion

We used both a modified APEC [2] code and an additional code written in Mathematica to solve equation (1). We calculated the ratio R for various conditions of the plasma heating and cooling.

It can be seen that in the case of fast heating of low density plasma the ratio R decreases dramatically for a short time (see the “valley” in Fig. 1b). A similar behavior holds for the case of cooling: we can see the dramatic increase of the ratio R in the first second (the “plateau” in Fig. 1d). For stationary case, R holds in the interval 1–10.

Unfortunately, such fast processes could not be observed, since the exposure time of X-ray satellites is of the order of 10^4 s and above. That is why we studied the influence of non-stationary processes on the average f/i -ratio (R_a). The model

parameters are presented in Table 1, results of the modeling are summarized in Table 2. It can be seen that the models E and F show significant difference between R_a and R_c and R_h . It means that using the ratios R for the stationary plasma leads to the overestimation of the n_e derived from observations.

The parameters of heating and cooling which are given in the captions to Fig. 1 are typical of solar nano-flares or similar events in stellar atmospheres.

4 Conclusion

We showed that non-stationary processes could affect on the instantaneous forbidden-to-intercombination ratio R_m which increases (for fast cooling) or decreases (for fast heating) by up to 1–3 orders of magnitude during the first second. These processes could also strongly change the averaged ratio R_a by decreasing it by up to 20%. In this case one can incorrectly estimate the plasma electron density (errors can be of up to 2 orders of magnitude), when supposing the stationary level population in the plasma.

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